Vortices in vibrated granular rods

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We report the experimental observation of vortex patterns in vertically vibrated granular rods. Above a critical packing fraction, moving ordered domains of nearly vertical rods spontaneously form and coexist with horizontal rods. The domains of vertical rods coarsen in time to form large vortices. We investigate the conditions under which the vortices occur by varying the number of rods, vibration amplitude, and frequency. The size of the vortices increases with the number of rods. We characterize the growth of the ordered domains by measuring the area fraction of the ordered regions as a function of time. A void-filling model is presented to describe the nucleation and growth of the vertical domains. We track the ends of the vertical rods and obtain the velocity fields of the vortices. The rotation speed of the rods is observed to depend on the vibration velocity of the container and on the packing. To investigate the impact of the direction of driving on the observed phenomena, we performed experiments with the container vibrated horizontally. Although vertical domains form, vortices are not observed. We therefore argue that the motion is generated due to the interaction of the inclination of the rods with the bottom of a vertically vibrated container. We also perform simple experiments with a single row of rods in an annulus. These experiments directly demonstrate that the rod motion is generated when the rods are inclined from the vertical, and is always in the direction of the inclination.

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I. INTRODUCTION

Granular materials are well known examples of dissipative nonequilibrium systems that show a rich variety of collective phenomenon, such as convection, wave patterns, and segregation [1–4]. Most studies utilize spherical particles to investigate these bulk properties. However, in most natural or industrial settings one can find an abundance of anisotropic granular materials, viz., rice, medicine capsules, and even logs. Therefore it is surprising that very few studies of prolate granular materials have been carried out. Mounfield and Edwards [5] applied the concepts of configurational statistical mechanics to study the nature of the isotropic to nematic phase transition in a granular system of elongated particles. In recent experiments utilizing a tall narrow cylinder, Villarruel et al. [6] studied the effects of anisotropy on granular packing. They observed the appearance of smectic states with the direction given by the container walls.

In thermal systems, particle anisotropy is known to produce ordered states. Examples are rod and plate shaped colloids and liquid crystals, which show orientational order and form nematic and smectic phases [7–11]. The ordering mechanism was found to be entropically driven, i.e., as thermodynamic equilibrium is reached, the process of entropy maximization leads to long range order. It is not obvious that such mechanisms carry over to granular systems because the thermal energy scale is very much smaller than the potential energy required for particle rearrangement, and energy has to be supplied actively to produce sustained motion. Therefore, an interesting question arises—does shape anisotropy lead to self-organization and pattern formation in granular materials?

In this paper we report the observation of vortex patterns exhibited by granular rods that are vibrated vertically inside a container. We obtain the phase diagram for the observed patterns as a function of the acceleration of the driving and the packing fraction of the rods. We find that for sufficiently large packing fractions, the rods tend to align vertically and undergo vortex motion. Using high frame rate imaging and particle tracking, we have measured the velocity fields of the vortices as a function of packing fraction and driving frequency. Based on our experimental observations, we argue that the inclination of the rods causes motion due to collisions with the bottom boundary. To bolster our claim, we conduct experiments with a single row of rods in an annulus to demonstrate that motion is always in the direction of inclination. Thus we show that shape anisotropy leads to translational motion in such systems.

II. EXPERIMENTAL APPARATUS

The experiments were performed in a circular anodized aluminum container with a diameter $D=6.0$ cm and depth $H=1.5$ cm. (Limited experiments were also carried out in a container with $D=8.5$ cm.) The container is leveled to within 0.002 cm and is attached to an electromechanical shaker through a rigid linear bearing that allows motion only in the vertical direction. The cell is driven by a sinusoidal signal at a frequency $f$, and is monitored by an accelerometer via a lock-in amplifier. The experiments were conducted using copper cylinders with uniform length $l=6.2$ mm and diameter $d=0.5$ mm. The cylindrical surface of the rods is coated with a gray tin-oxide layer to diminish light reflection and has a coefficient of friction $\mu=0.32$. The patterns that form are imaged from above using a high-frame rate digital camera (Kodak SR-1000). Since the flat ends of the rods reflect light better than the sides, they appear as bright spots when imaged from above. If the rods are inclined greater than $35^\circ$, they reflect far less light compared to nearly vertical rods. If the rods are inclined greater than $35^\circ$, they reflect far less light compared to nearly vertical rods. If the rods are inclined greater than $35^\circ$, they reflect far less light compared to nearly vertical rods.

One control parameter for our experiments is $\Gamma = A(2\pi f)^2/g$, where $A$ is the driving amplitude and $g$ is the acceleration due to gravity.
acceleration due to gravity. A second control parameter is the non-dimensional number fraction \( n_\phi = N / N_{\text{max}} \), where \( N \) is the total number of rods in the container and \( N_{\text{max}} = (\pi / \sqrt{12})(D / d)^2 \) is the maximum number of rods required to obtain a vertically aligned monolayer. Therefore, \( n_\phi = 1 \) corresponds to having one layer of triangularly packed vertical rods.

### III. OBSERVED PATTERNS

The system is initialized by pouring the rods into the container and then increasing \( \Gamma \), which establishes a random state. Figure 1 shows the phase diagram of the observed patterns as a function of \( n_\phi (f = 50 \text{ Hz}) \). The rods are observed to remain static in the initial configuration until \( \Gamma \geq 1.47 \). As \( \Gamma \) is increased above this value, the rods heap towards one side of the container similar to previous observations with spherical particles [12]. As the acceleration is increased further above \( \Gamma = 2.6 \), the rods are observed to spread out evenly inside the container, and form horizontal layers with random orientations. An example of such a nematic phase observed at low \( n_\phi \) is shown in Fig. 2(a). Above a critical \( n_\phi \), a second transition is observed. Chaotically moving domains of almost vertical rods are observed to spontaneously form and coexist with horizontal rods [13]. An example is shown in Fig. 2(b).

As \( n_\phi \) is increased further, the domains of near-vertical rods that are formed coalesce and undergo vortex motion. Figure 2(c) shows an example of two counterrotating vortices. For large number fractions \( (n_\phi > 0.49) \), a single large stable vortex made entirely of near-vertical rods is observed near the center and inclined rods at its boundary is observed [Fig. 2(d)].

If the acceleration is increased further, a gas-like state is reached where the rods vibrate vigorously inside the container and the nematic and vortex states are destroyed. For the highest \( n_\phi \), a pure gaseous state is not reached due to the limits of the apparatus. Instead, domains of ordered vertical rods are observed to coexist with the gas of rods.

![Phase diagram](image1)

**FIG. 1.** (Color online) The phase diagram of the observed patterns in vibrated granular rods. The driving frequency \( f = 50 \text{ Hz} \). Vortices are observed for sufficiently large numbers of rods and driving acceleration parameter \( \Gamma \).**

**FIG. 2.** Examples of patterns observed when granular rods are vertically vibrated \((f = 50 \text{ Hz})\). (a) Nematiclike state \((n_\phi = 0.152, \Gamma = 2.42)\). (b) Moving domains of near vertical rods \((n_\phi = 0.344, \Gamma = 3.38)\). (c) Multiple vortices \((n_\phi = 0.551, \Gamma = 3.29)\). (d) Large vortex \((n_\phi = 0.535, \Gamma = 3.00)\). The cylindrical surface of the rods is coated with a gray tin-oxide layer to diminish light reflection and has a coefficient of friction \( \mu = 0.32 \). The patterns that form are imaged from above using a high-frame rate digital camera (Kodak SR-1000). Since the tips of the rods reflect light better than the sides, they appear as bright spots when imaged from above. If the rods are inclined greater than 35°, they reflect far less light compared to nearly vertical and horizontal rods.

#### A. Domain growth

We now discuss how vortices nucleate and grow as a function of time. An example of the growth process is shown in Fig. 3, and a movie can be viewed at Ref. [14]. Starting from a random state, pockets of near-vertical domains are observed to nucleate uniformly inside the container. This is in contrast with previous observations where the ordered domains grow inward from the boundary [6]. The pockets of vertical domains grow in time by merging to form larger domains of almost vertical rods. The domains are then observed to collectively move as the system becomes more ordered and then finally show vorticity.

We measure the growth of the domains by taking time-lapse images of the process described above. The growth of near-vertical domains is then measured from the ratio of the domain area to that of the total container area as a function of time. The fraction of rods that are near-vertical, \( n_v \), is plotted in Fig. 4 for a range of \( n_\phi \) at a fixed driving acceleration and frequency, \( \Gamma = 3.4 \) and \( f = 50 \text{ Hz} \).

We find that there exist two growth regimes (a model is discussed in Sec. V). For short times, the growth rate depends on \( n_\phi \) in the following way. If \( n_\phi \) is below a critical value, domains appear but never organize into large vortices. However, as \( n_\phi \) is increased the short-time growth is slowed by the high packing. That is, more time is required for small domains to form due to the lack of voids present. A void
filling model for domain growth will be discussed in Sec. V. The saturation value of $n_v$ increases with $n_f$. The difference in the values arises due to the definitions of $n_v$ and $n_f$.

The variation of the relative number of near-vertical rods as a function of $n_f$ is plotted in Fig. 5. We show separate data for the relative number of inclined rods, vertical rods and horizontal rods. Rods that are tilted between 30° and 80° are considered inclined, and the rest are considered as either horizontal or vertical (the inclination of the rod was obtained by measuring the distance $x$ between adjacent rods and calculating $\sin^{-1}(d/x)$). For $n_f<0.22$, the container has only horizontal rods. As $n_f$ increases, more rods tend to the vertical direction. For $n_f>0.71$ most rods are either vertical or horizontal.

B. Spatial structure of the vortex and velocity fields

We next discuss the spatial structure of the vortex. In Fig. 6 a close-up image of a vortex pattern is displayed which shows the progressive inclination of the rods from the center of the vortex. We note that the direction of motion corresponds to the direction of inclination (for Fig. 6 the motion is counterclockwise). In this case the inclination of the rods varies from approximately 85° to 40° near the edges. The
change in inclination of the rods within the vortex is less at higher \( n_f \) due to the increased packing.

By acquiring images at 250 (frames s\(^{-1}\)), we were able to track the ends of the vertical rods and obtain the velocity fields of the vortices. In addition to the rotational motion, rods also vibrated due to the collisional interactions with their neighbors. We first track \( \partial_1 \) the ends of the near-vertical rods and then perform spatial and temporal averaging to obtain the vortex fields. An example of the velocity field of a vortex is presented in Fig. 7 (see text).

The averaged azimuthal velocity \( v(r) \), as a function of the distance \( r \) from the center of the vortex, is shown in Fig. 8\( \text{a} \) for a range of \( n_f \). The data was obtained at a constant acceleration \( \Gamma = 3.00 \) for a frequency \( f = 50 \) Hz. We note that for small distances relative to the center of the vortices, the averaged velocity \( v(r) \) increases linearly with the distance \( r \) indicating solid body rotation. At intermediate ranges, this linear relationship is not observed, and therefore shearing occurs inside the vortices. As the boundary of the vortex is approached, velocity decreases due to friction with the horizontal rods at the edge. It is interesting to note that the horizontal rods at the edge of the vortex are often aligned tangentially with the boundary of the vortex [see Fig. 2(c)]. We also observe that the slope of the averaged velocity \( v(r) \) at small \( r \) systematically decreases as the packing fraction \( n_\phi \) is increased. Thus the angular velocity of the inner core of the vortex decreases with its size.

**C. Frequency dependence**

We also explored the rotation rate dependence of the vortices as a function of \( f \). Figure 8\( \text{b} \) shows \( v(r) \) for \( n_\phi = 0.53 \) and \( \Gamma = 3.00 \) for vortices of equivalent size. As stated, \( \Gamma \) is held constant while frequency is increased, which decreases vibration velocity. Therefore, we demonstrate that the vortex speed systematically increases with vibration velocity.

**IV. HORIZONTALLY VIBRATED CONTAINER**

To test if the transition to a vertical state is dependent on the way the container is vibrated, experiments were performed with horizontal driving. There do in fact exist similarities: at low \( n_\phi \) heaping and nematiclike domains of horizontal rods are observed; at high \( n_\phi \) the rods tend to align vertically and form ordered domains. An image is shown in Fig. 9. In contrast to the vertical shaking the vertically aligned domains do not migrate and coarsen into vortices. We also find that in this case the convection is stronger and destroys the motion of the vertically ordered domains.
V. DISCUSSION

We first discuss possible mechanisms responsible for the
rod to align vertically at high packing fractions. The ten-
dency of the rods to align vertically at high packing fractions
may be understood in terms of a void filling mechanism. If
voids (i.e., space between rods) are small then they can only
be filled with a near-vertical rod. On the other hand, larger
voids will accommodate a horizontal rod. Assuming that the
distribution of voids decreases with void size, the most prob-
able configuration will be regions of vertical rods. Further-
more, the decrease in the number of large voids is enhanced
by the coalescence of regions of vertically aligned rods. This
also drives the system to pack more closely and produces a
decreased overall center of mass. Thus ordering at high pack-
ing fractions can be explained by a process that is analogous
to configurational entropy driven ordering in thermal systems
[5].

To describe the nucleation and growth of the vertical do-
main at high \( n_\phi \) shown in Figs. 3 and 4 we present a simple
analysis. The rods are assumed to be either vertical or hori-
zontal. We then assume that the growth rate is initially linear
and asymptotically must decrease to zero as the number of
horizontal rods is diminished and a steady state is reached.
Then the evolution of \( n_v \) is described by

\[
\frac{\partial n_v}{\partial t} = \alpha n_v(\beta - n_v),
\]

where \( \alpha \) and \( \beta \) are constants that depend on \( n_\phi \). Equation
(1) can be integrated and is fit to the measured \( n_v(t) \) in Fig.
4. This simplified interpretation seems to capture the nucle-
ation, growth, and saturation of the near-vertical domains in
our experiments.

In previous work by Villarruel et al. [6] the experiments
were carried out in a system where the diameter of the con-
tainer and the length of the rods were comparable. The con-
tainer was tapped and thus the rods experienced considerable
shearing with the side boundary. They observed that vertical
domains nucleated at the boundaries and subsequently propa-
gate to the center of the cell. These observations are consis-
tent with our findings. We have shown that vertical domains
can nucleate away from the sidewalls and form independent
of the direction of the driving.

Next, we elucidate the physical mechanism that is respon-
sible for the vortex motion. We performed additional experi-
ments with a row of cylindrical rods with \( l = 5.1 \text{ cm} \) and \( d = 0.6 \text{ cm} \)
in a one-dimensional 1D annulus with a mean ra-
radius of 5.5 cm (see Fig. 10). When subjected to vertical
motion, the rods were always observed to move in the direc-
tion of the inclination. No translation motion is observed
when the rods were vertical. A movie showing this property
can be found at Ref. [14]. We also used ball point pens and
pencils to check that the detailed shape of the rod and the tip
is not important to this translation mechanism. We find that
the translation speed depends on the driving frequency and
the inclination is qualitatively similar to that observed for the
vortices.

The physical mechanism for the motion of the inclined
rods, based on our observation, appears to be as follows.
When the inclined rods are vibrated vertically, they hit the
bottom plate at a point away from their center of mass. Be-
cause the rotation of the rod about its center of mass is con-
strained due to the neighboring rods, it gets launched in the
direction of the inclination. Thus, the greater the inclination
of the rods, the further they get launched and land (viz.,
projectile motion), giving rise to translational motion.

From these direct observations we conclude that the in-
clined rods form the engine that drives the vortex. The ver-
tical rods in the center of the vortex are simply pulled around
by shear induced by the inclined rods. As the number of rods
is increased and the size of the vortex increases, the inclina-
tion of the rods also decreases because of greater packing.
Therefore, the speed of the very large vortices is expected to
decrease with the size of the vortex, consistent with our ob-
servation [see Fig. 8(a)]. We also note that the frequency
dependence of \( \nu(r) \) is compatible with the mechanism de-
scribed above. By decreasing the frequency of the driving signal at fixed $\Gamma$ the velocity of the driving plate is increased. Therefore, one expects $v(r)$ to increase with a decrease in frequency, consistent with our observations [see Fig. 8(b)].

VI. CONCLUSION

In conclusion, granular rods are observed to form vertically aligned domains at high packing fractions, independent of the direction of vibration. Vortex patterns are observed when the rods are vibrated vertically. We have measured the growth of near-vertical domains and have developed a simple void filling model that well describes our results. We have also shown a new translation mechanism which occurs due to anisotropy. Based on these observations, Aranson and Tsimring [16] have developed a phenomenological model that describes the formation and coarsening of the vortices.

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[14] See EPAPS Document No. E-PLEEE8-67-161303 for movies showing the vorticity of the patterns (also available at http://physics.clarku.edu/vortex). A direct link to this document may be found in the online article’s HTML reference section. The document may also be reached via the EPAPS homepage (http://www.aip.org/pubservs/epaps.html) or from ftp.aip.org in the directory /epaps/. See the EPAPS homepage for more information.